

**BOUNDEDNESS OF LARGE-TIME SOLUTIONS  
TO A CHEMOTAXIS MODEL  
WITH NONLOCAL AND SEMILINEAR FLUX**

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**ABSTRACT.** A semilinear version of parabolic-elliptic Keller–Segel system with the *critical* nonlocal diffusion is considered in one space dimension. We show boundedness of weak solutions under very general conditions on our semilinearity. It can degenerate, but has to provide a stronger dissipation for large values of a solution than in the critical linear case or we need to assume certain (explicit) data smallness. Moreover, when one considers a logistic term with a parameter  $r$ , we obtain our results even for diffusions slightly weaker than the critical linear one and for arbitrarily large initial datum, provided  $r > 1$ . For a mild logistic dampening, we can improve the smallness condition on the initial datum up to  $\sim 1/(1 - r)$ .

### 1. Introduction

In this paper we study the following model:

$$(1.1) \quad \partial_t u = \partial_x(-\mu(u)Hu + u\partial_x v) + ru(1 - u), \quad x \in \mathbb{T}, t \in \mathbb{R}^+,$$

$$(1.2) \quad \partial_x^2 v = u - \langle u \rangle, \quad x \in \mathbb{T}, t \in \mathbb{R}^+,$$

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where  $u = u(x, t)$ ,  $v = v(x, t)$ ,  $H$  stands for the (periodic) Hilbert transform, i.e.

$$\widehat{Hu}(\xi) = -i \frac{\xi}{|\xi|} \widehat{u}(\xi).$$

$\mathbb{T} = [-\pi, \pi]$ ,  $r \geq 0$  and  $\mu$  is a certain function (semilinearity), precised in what follows. Before formulating our results let us explain our motivations to study system (1.1)–(1.2).

**1.1. Motivation.** (a) *Mathematical biology.* One of the basic systems studied in the context of chemotaxis is the parabolic–elliptic Keller–Segel system (also known as the Smoluchowski–Poisson system)

$$(1.3) \quad \partial_t u = \nabla \cdot (\mu \nabla u - u \nabla \phi), \quad x \in \mathbb{T}, \quad t \in \mathbb{R}^+,$$

where  $d \geq 1$  denotes the spatial dimension,  $\mathbb{T}^d = [-\pi, \pi]^d$ ,  $\mu > 0$  is a constant and  $\phi$  is recovered from  $u$  through some operator, i.e.  $\phi(x, t) = T(u(x, t))$ . In many cases  $\phi$  satisfies the Poisson equation

$$(1.4) \quad -\Delta \phi = u - \langle u \rangle, \quad x \in \mathbb{T}^d, \quad t \in \mathbb{R}^+.$$

In this notation,  $u$  represents the concentration of cells,  $\langle u \rangle$  its space average and  $\phi$  gives us the concentration of a chemical substance that attracts cells. It is biologically justified to enrich equation (1.3) with the logistic term, obtaining

$$(1.5) \quad \partial_t u = \nabla \cdot (\mu \nabla u - u \nabla \phi) + ru(1 - u), \quad x \in \mathbb{T}^d, \quad t \in \mathbb{R}^+,$$

where  $r \geq 0$ . Model (1.5)–(1.4) is related to the parabolic-elliptic simplification of the cell kinetics model M8 in [28], that describes a bacterial pattern formation or cell movement and growth during angiogenesis.

Another application of model (1.5)–(1.4) occurs in tumor growth. In particular, this model is related to the three-component urokinase plasminogen invasion model (see [29]). There is a huge literature on the mathematical study of a numerous versions of (1.5)–(1.4) in the context of mathematical biology, see [5], [7], [9], [10], [16], [25], [30] and the references therein.

(b) *Natural sciences.* Let us take in (1.5)–(1.4),  $v := -\phi$ . The resulting system

$$(1.6) \quad \partial_t u = \nabla \cdot (\mu \nabla u + u \nabla v) + ru(1 - u), \quad x \in \mathbb{T}^d, \quad t \in \mathbb{R}^+,$$

$$(1.7) \quad \Delta v = u - \langle u \rangle, \quad x \in \mathbb{T}^d, \quad t \in \mathbb{R}^+,$$

in the case  $r = 0$  is important in mathematical cosmology and gravitation theory. It is very similar in spirit to the Zel'dovich approximation used in cosmology to study the formation of large-scale structure in the primordial universe, see also [1], [4]. It is also connected with the Chandrasekhar equation for the gravitational equilibrium of polytropic stars, statistical mechanics and the Debye system for electrolytes, see [6].